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ON A BOUNDARY VALUE PROBLEM FOR A FOURTH ORDER PARTIAL DIFFERENTIAL EQUATION

Abstract

The existence of a classic solution of a boundary value problem for a fourth order partial differential equation is proved.

Recently, there is a great interest to the problems of wave propagation in stratified liquids in connection with needs of oceanology and applied geophysics.

In the paper [1] it was derived the main equation of dynamics of compressible exponentially stratified liquid:

$$\frac{\partial^2}{\partial t^2} \left[\frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} - \Delta_3 u + \beta^2 u \right] - \omega_0^2 \Delta_2 u = 0, \tag{*}$$

here Δ_3 and Δ_2 are Laplace operators with respect to variables (x_1, x_2, x_3) and (x_1, x_2) , respectively. The fourth order differential equation

$$\frac{\partial^2}{\partial t^2} \left[\frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} - u_{xx} + \beta^2 u \right] - \omega_0^2 u_{xx} = 0, \tag{**}$$

being one-dimensional analogy of general equation (*) was considered in [2]. Parameter c denotes sound velocity in the considered compressible exponentially stratified liquid; the quantity 2β equals index in the exponent characterizing density distribution of liquid and in accepted approximation may be considered to be constant; main frequency parameter of the considered liquid called Weysel-Brandtl frequency is denoted by ω_0 . Its quantity is determined by the equality

$$\omega_0^2 = 2\beta g - g^2/c^2,$$

where g is free fall acceleration.

For convenience of further considerations we pass to non-dimensional independent variables $\beta x, \omega_0 t$ and keep previous denotation x and t . The equation (**) takes the form:

$$\varepsilon^2 \frac{\partial^4 u}{\partial t^4} - \frac{\partial^2}{\partial t^2} [u_{xx} - u] - u_{xx} = 0 \tag{***}$$

where $\varepsilon = \omega_0 / (\beta c) \leq 1$. Fundamental solution of equation (***) is studied in [2].

Now, let's consider the boundary-value problem:

$$u_{tttt}(x, t) - u_{ttxx}(x, t) + u_{tt}(x, t) - u_{xx}(x, t) = F(x, t, u(x, t), u_t(x, t)),$$

$$(x, t) \in D_T \equiv \{(x, t) : 0 \leq x \leq 1, 0 \leq t \leq T\}, \tag{1}$$

$$u(0, t) = 0, \quad u_x(1, t) = 0, \quad 0 \leq t \leq T, \tag{2}$$

$$u(x, 0) = \varphi_0(x), \quad u_t(x, 0) = \varphi_1(x),$$

$$u_{tt}(x, 0) = \varphi_2(x), \quad u_{ttt}(x, 0) = \varphi_3(x), \quad 0 \leq x \leq 1, \quad (3)$$

where

$$F(x, t, u(x, t), u_t(x, t)) = p_0(t)u(x, t) + p_1(t)u_t(x, t) + f(x, t), \quad (4)$$

$f(x, t), \varphi_0(x), \varphi_1(x), \varphi_2(x), \varphi_3(x), p_0(t), p_1(t)$ are the given functions, $u(x, t)$ is the desired function.

Under the classic solution of problem (1)-(3) we understand the functions $u(x, t)$, continuous in closed domain D_T together with its derivatives contained in equation (1) and satisfying all conditions of (1)-(3) in an ordinary sense.

By $B_{2,T}^{\alpha,\beta}$ ([3]) we denote aggregate of all functions of the form

$$u(x, t) = \sum_{k=1}^{\infty} u_k(t) \sin \lambda_k x, \quad \lambda_k = \frac{\pi}{2}(2k - 1),$$

considered in D_T where each of the functions $u_k(t)$ is continuously differentiable on $[0, T]$, and

$$J(u) \equiv \left\{ \sum_{k=1}^{\infty} \left(\lambda_k^\alpha \|u_k(t)\|_{C[0,T]} \right)^2 \right\}^{\frac{1}{2}} + \left\{ \sum_{k=1}^{\infty} \left(\lambda_k^\beta \|u'_k(t)\|_{C[0,T]} \right)^2 \right\}^{\frac{1}{2}} < +\infty$$

moreover $\alpha \geq 0, \beta \geq 0$.

In this set we define the norm as follows:

$$\|u(x, t)\|_{B_{2,T}^{\alpha,\beta}} = J(u).$$

It is known that [3] $B_{2,T}^{\alpha,\beta}$ are Banach spaces.

Since the system $\{\sin \lambda_k x\}_{k=1}^{\infty}$ is complete in $L_2(0, 1)$, obviously each classic solution $u(x, t)$ of problem (1)-(3) is of the form:

$$u(x, t) = \sum_{k=1}^{\infty} u_k(t) \sin \lambda_k x, \quad (5)$$

where

$$u_k(t) = 2 \int_0^1 u(x, t) \sin \lambda_k x dx.$$

Now we apply Fourier method and get from (1)-(3):

$$u_k^{(4)}(t) + (1 + \lambda_k^2) u_k''(t) + \lambda_k^2 u_k(t) = F_k(t; u), \quad (6)$$

$$u_k(0) = \varphi_{0k}, \quad u'_k(0) = \varphi_{1k}, \quad u''_k(0) = \varphi_{2k}, \quad u_k'''(0) = \varphi_{3k} \quad (7)$$

where

$$F_k(t; u) \equiv p_0(t)u_k(t) + p_1(t)u'_k(t) + f_k(t),$$

$$f_k(t) = 2 \int_0^1 f(x, t) \sin \lambda_k x dx,$$

$$\varphi_{ik} = 2 \int_0^1 \varphi_i(x) \sin \lambda_k x dx \quad (i = \overline{0, 3}) \quad (8)$$

The roots of the characteristic function

$$\mu_{4k} + (1 + \lambda_k^2) \mu_k^2 + \lambda_k^2 = 0$$

corresponding to (6) are determined by the relations:

$$\begin{aligned} \mu_{jk} &= (-1)^j \sqrt{-1} \quad (j = 1, 2), \\ \mu_{jk} &= (-1)^j \lambda_k \sqrt{-1} \quad (j = 3, 4). \end{aligned}$$

Then, after applying the method of variation of parameters we reduce the solution of problem (6), (7) to the solution of the following denumerable system:

$$\begin{aligned} u_k(t) &= \frac{1}{\lambda_k^2 - 1} \left[(\lambda_k^2 \cos t - \cos \lambda_k t) \varphi_{0k} + \left(\lambda_k^2 \sin t - \frac{\sin \lambda_k t}{\lambda_k} \right) \varphi_{1k} + \right. \\ &\quad + (\cos t - \cos \lambda_k t) \varphi_{2k} + \left(\sin t - \frac{\sin \lambda_k t}{\lambda_k} \right) \varphi_{3k} + \\ &\quad \left. + \frac{1}{\lambda_k} \int_0^t F_k(\tau; u) (\lambda_k \sin(t - \tau) - \sin \lambda_k(t - \tau)) d\tau \right]. \quad (9) \end{aligned}$$

Substituting $u_k(t)$ form (9) to the representation of function (5) we get:

$$\begin{aligned} u(x, t) &= \sum_{k=1}^{\infty} \left\{ \frac{1}{\lambda_k^2 - 1} \left[(\lambda_k^2 \cos t - \cos \lambda_k t) \varphi_{0k} + \left(\lambda_k^2 \sin t - \frac{\sin \lambda_k t}{\lambda_k} \right) \varphi_{1k} + \right. \right. \\ &\quad + (\cos t - \cos \lambda_k t) \varphi_{2k} + \left(\sin t - \frac{\sin \lambda_k t}{\lambda_k} \right) \varphi_{3k} + \\ &\quad \left. \left. + \frac{1}{\lambda_k} \int_0^t F_k(\tau; u) (\lambda_k \sin(t - \tau) - \sin \lambda_k(t - \tau)) d\tau \right] \right\} \sin \lambda_k t. \quad (10) \end{aligned}$$

From (9) we have:

$$\begin{aligned} u'_k(t) &= \frac{1}{\lambda_k^2 - 1} \left[(-\lambda_k^2 \sin t + \lambda_k \sin \lambda_k t) \varphi_{0k} + (\lambda_k^2 \cos t - \cos \lambda_k t) \varphi_{1k} + \right. \\ &\quad + (-\sin t + \lambda_k \sin \lambda_k t) \varphi_{2k} + (\cos t - \cos \lambda_k t) \varphi_{3k} + \\ &\quad \left. + \int_0^t F_k(\tau; u) (\cos(t - \tau) - \cos \lambda_k(t - \tau)) d\tau \right], \quad (11) \end{aligned}$$

$$\begin{aligned} u''_k(t) &= \frac{1}{\lambda_k^2 - 1} \left[(-\lambda_k^2 \cos t + \lambda_k^2 \cos \lambda_k t) \varphi_{0k} + (-\lambda_k^2 \sin t + \lambda_k \sin \lambda_k t) \varphi_{1k} + \right. \\ &\quad \left. + (-\cos t + \lambda_k^2 \sin \lambda_k t) \varphi_{2k} + (-\sin t + \lambda_k \sin \lambda_k t) \varphi_{3k} + \right. \end{aligned}$$

$$+ \int_0^t F_k(\tau; u) (-\sin(t-\tau) + \lambda_k \sin \lambda_k(t-\tau)) d\tau \Big], \quad (12)$$

$$\begin{aligned} u_k'''(t) & \frac{1}{\lambda_k^2 - 1} [(\lambda_k^2 \sin t - \lambda_k^3 \sin \lambda_k t) \varphi_{0k} + (\lambda_k^2 \cos t + \lambda_k^2 \cos \lambda_k t) \varphi_{1k} + \\ & + (\sin t - \lambda_k^3 \sin \lambda_k t) \varphi_{2k} + (-\cos t + \lambda_k^2 \cos \lambda_k(t-\tau)) \varphi_{3k} + \\ & + \int_0^t F_k(\tau; u) (-\cos(t-\tau) + \lambda_k^2 \cos \lambda_k(t-\tau)) d\tau \Big], \quad (13) \end{aligned}$$

$$\begin{aligned} u_k^4(t) & = \frac{1}{\lambda_k^2 - 1} [(\lambda_k^2 \cos t - \lambda_k^4 \cos \lambda_k t) \varphi_{0k} + (\lambda_k^2 \sin t - \lambda_k^3 \sin \lambda_k t) \varphi_{1k} + \\ & + (\cos t - \lambda_k^4 \cos \lambda_k t) \varphi_{2k} + (\sin t - \lambda_k^3 \sin \lambda_k t) \varphi_{3k} + \\ & + \int_0^t F_k(\tau; u) (\sin(t-\tau) - \lambda_k^3 \sin \lambda_k(t-\tau)) d\tau \Big] + F_k(t; u). \quad (14) \end{aligned}$$

Proceeding from the determination of a classic solution of problem (1)-(3) we easily solve the following

Lemma. *If $u(x, t)$ is any classic solution of problem (1)-(3), the functions*

$$u_k(t) = 2 \int_0^1 u(x, t) \sin \lambda_k x dx \quad (k = 1, 2, \dots)$$

satisfy system (9) on $[0, T]$.

We have from (9), (11)-(14):

$$\begin{aligned} |u_k(t)| & \leq \frac{2}{\lambda_k^2} (1 + \lambda_k^2) |\varphi_{0k}| + \frac{2}{\lambda_k^3} (1 + \lambda_k^3) |\varphi_{1k}| + \frac{4}{\lambda_k^2} |\varphi_{2k}| + \\ & + \frac{2}{\lambda_k^3} (1 + \lambda_k) |\varphi_{3k}| + \frac{2}{\lambda_k^3} (1 + \lambda_k) \sqrt{T} \left(\int_0^T |F_k(\tau; u)|^2 d\tau \right)^{\frac{1}{2}}, \\ |u_k'(t)| & \leq \frac{2}{\lambda_k^2} (\lambda_k + \lambda_k^2) |\varphi_{0k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k^2) |\varphi_{1k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k) |\varphi_{2k}| + \\ & + \frac{4}{\lambda_k^2} |\varphi_{3k}| + \frac{4}{\lambda_k^2} \sqrt{T} \left(\int_0^T |F_k(\tau; u)|^2 d\tau \right)^{\frac{1}{2}}, \\ |u_k''(t)| & \leq 4 |\varphi_{0k}| + \frac{2}{\lambda_k^2} (\lambda_k + \lambda_k^2) |\varphi_{1k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k^2) |\varphi_{2k}| + \\ & + \frac{2}{\lambda_k^2} (1 + \lambda_k) |\varphi_{3k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k) \sqrt{T} \left(\int_0^T |F_k(\tau; u)|^2 d\tau \right)^{\frac{1}{2}}, \end{aligned}$$

$$\begin{aligned}
 |u_k'''(t)| &\leq 2(1 + \lambda_k) |\varphi_{0k}| + 4|\varphi_{1k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k) |\varphi_{2k}| + \\
 &+ \frac{2}{\lambda_k^2} (1 + \lambda_k) |\varphi_{3k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k) \sqrt{T} \left(\int_0^T |F_k(\tau; u)|^2 d\tau \right)^{\frac{1}{2}}, \\
 |u_k^{(4)}(t)| &\leq 2(1 + \lambda_k^2) |\varphi_{0k}| + 2(1 + \lambda_k) |\varphi_{1k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k^4) |\varphi_{2k}| + \\
 &+ \frac{2}{\lambda_k^2} (1 + \lambda_k^3) |\varphi_{3k}| + \frac{2}{\lambda_k^2} (1 + \lambda_k) \sqrt{T} \left(\int_0^T |F_k(\tau; u)|^2 d\tau \right)^{\frac{1}{2}} + |F_k(t; u)|.
 \end{aligned}$$

Hence we have:

$$\begin{aligned}
 &\left(\sum_{k=1}^{\infty} \left(\lambda_k^3 \|u_k(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} \leq \\
 &\leq 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{0k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{1k}|)^2 \right)^{\frac{1}{2}} + \\
 &+ 4 \left(\sum_{k=1}^{\infty} (\lambda_k |\varphi_{2k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k |\varphi_{3k}|)^2 \right)^{\frac{1}{2}} + \\
 &+ 4T \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} + \\
 &+ 4T \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k'(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} + \\
 &+ 4\sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}}; \tag{15} \\
 &\left(\sum_{k=1}^{\infty} \left(\lambda_k^3 \|u_k'(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} \leq \\
 &\leq 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{0k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{1k}|)^2 \right)^{\frac{1}{2}} + \\
 &+ 4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{2k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k |\varphi_{3k}|)^2 \right)^{\frac{1}{2}} + \\
 &+ 4T \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} +
 \end{aligned}$$

$$\begin{aligned}
& +4T \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k \|u'_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\
& +4\sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}} ; \tag{16}
\end{aligned}$$

$$\begin{aligned}
& \left(\sum_{k=1}^{\infty} (\lambda_k^3 \|u''_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \leq 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{0k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{1k}|)^2 \right)^{\frac{1}{2}} + \\
& +4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{2k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{3k}|)^2 \right)^{\frac{1}{2}} + \\
& +4T \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k^2 \|u'_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\
& +4T \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k^2 \|u'_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\
& +4\sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k^2 |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}} . \tag{17}
\end{aligned}$$

$$\begin{aligned}
& \left(\sum_{k=1}^{\infty} (\lambda_k \|u'''_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \leq \\
& \leq 4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{0k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k |\varphi_{1k}|)^2 \right)^{\frac{1}{2}} + \\
& +4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{2k}|)^2 \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k |\varphi_{3k}|)^2 \right)^{\frac{1}{2}} + \\
& +4T \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\
& +4T \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k \|u'_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\
& +4\sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}} . \tag{18}
\end{aligned}$$

$$\left(\sum_{k=1}^{\infty} (\lambda_k \|u_k^{(4)}(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \leq$$

$$\begin{aligned}
 &\leq 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{0k}|) \right)^2 + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{1k}|) \right)^{\frac{1}{2}} + \\
 &+ 4 \left(\sum_{k=1}^{\infty} (\lambda_k^3 |\varphi_{2k}|) \right)^{\frac{1}{2}} + 4 \left(\sum_{k=1}^{\infty} (\lambda_k^2 |\varphi_{3k}|) \right)^{\frac{1}{2}} + \\
 &+ 4T \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k^2 \|u_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} + \\
 &+ 4T \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k^2 \|u'_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} + 4\sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} \lambda_k^2 |f_k(\tau)|^2 d\tau \right)^{\frac{1}{2}} + \\
 &+ \|p_0(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k \|u_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} + \\
 &+ \|p_1(t)\|_{C[0,T]} \left(\sum_{k=1}^{\infty} (\lambda_k \|u_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} + \left(\sum_{k=1}^{\infty} (\lambda_k \|f_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}}. \quad (19)
 \end{aligned}$$

Assume that the problem's data satisfy the conditions:

1. $\varphi_i(x) \in C^2[0, 1]$, $\varphi_i'''(x) \in L_2(0, 1)$ and $\varphi_i(0) = \varphi_i'(1) = \varphi_i''(0) = 0$ ($i = 0, 1, 2$).
2. $\varphi_3(x) \in C^1[0, 1]$, $\varphi_3'''(x) \in L_2(0, 1)$ and $\varphi_3(0) = \varphi_3'(1) = 0$.
3. $f(x, t) \in C_{x,t}^{1,0}(D_T)$, $f_{xx}(x, t) \in L_2(D_T)$ and $f(0, t) = f_x(1, t) = 0$ ($0 \leq t \leq T$), ..
4. $p_0(t), p_1(t) \in C[0, T]$.

Then from (15)-(19) we get:

$$\begin{aligned}
 &\left(\sum_{k=1}^{\infty} (\lambda_k^3 \|u_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} \leq 4\sqrt{2} \|\varphi_0'''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_1'''(x)\|_{L_2(0,1)} + \\
 &+ 4\sqrt{2} \|\varphi_2''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_3'(x)\|_{L_2(0,1)} + 4\sqrt{2T} \|f_x(x, t)\|_{L_2(D_T)} + \\
 &+ 4T \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right) \|u\|_{B_{2,T}^{3,3}}; \quad (20)
 \end{aligned}$$

$$\begin{aligned}
 &\left(\sum_{k=1}^{\infty} (\lambda_k^3 \|u'_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} \leq 4\sqrt{2} \|\varphi_0'''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_1'''(x)\|_{L_2(0,1)} + \\
 &+ 4\sqrt{2} \|\varphi_2''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_3'(x)\|_{L_2(0,1)} + 4\sqrt{2T} \|f_x(x, t)\|_{L_2(D_T)} + \\
 &+ 4T \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right) \|u\|_{B_{2,T}^{3,3}}; \quad (21)
 \end{aligned}$$

$$\left(\sum_{k=1}^{\infty} (\lambda_k^3 \|u''_k(t)\|_{C[0,T]}) \right)^{\frac{1}{2}} \leq 4\sqrt{2} \|\varphi_0'''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_1'''(x)\|_{L_2(0,1)} +$$

$$+4\sqrt{2} \|\varphi_2'''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_3''(x)\|_{L_2(0,1)} + 4\sqrt{2T} \|f_{xx}(x,t)\|_{L_2(D_T)} + \\ +4T \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right) \|u\|_{B_{2,T}^{3,3}}; \quad (22)$$

$$\left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k'''(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} \leq 4\sqrt{2} \|\varphi_0''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_1'(x)\|_{L_2(0,1)} + \\ +4\sqrt{2} \|\varphi_2''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_3'(x)\|_{L_2(0,1)} + 4\sqrt{2T} \|f_x(x,t)\|_{L_2(D_T)} + \\ +4T \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right) \|u\|_{B_{2,T}^{3,3}}; \quad (23)$$

$$\left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k^{(4)}(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} \leq 4\sqrt{2} \|\varphi_0'''(x)\|_{L_2(0,1)} + \\ +4\sqrt{2} \|\varphi_1''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_2'''(x)\|_{L_2(0,1)} + 4\sqrt{2} \|\varphi_3''(x)\|_{L_2(0,1)} + \\ + \left(4\sqrt{T} + 1 \right) \sqrt{2} \|f_{xx}(x,t)\|_{L_2(D_T)} + \\ + (4T + 1) \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right) \|u\|_{B_{2,T}^{3,3}}; \quad (24)$$

Now we deduce from inequalities (20) and (21):

$$\|u\|_{B_{2,T}^{3,3}} \equiv \left(\sum_{k=1}^{\infty} \left(\lambda_k^3 \|u_k(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} + \\ + \left(\sum_{k=1}^{\infty} \left(\lambda_k^3 \|u_k'(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}} \leq A_1(T) + A_2(T) T \|u\|_{B_{2,T}^{3,3}}, \quad (25)$$

where

$$A_1(T) = 8\sqrt{2} \|\varphi_0'''(x)\|_{L_2(0,1)} + 8\sqrt{2} \|\varphi_1'''(x)\|_{L_2(0,1)} + 8\sqrt{2} \|\varphi_2''(x)\|_{L_2(0,1)} + \\ + 8\sqrt{2} \|\varphi_3'(x)\|_{L_2(0,1)} + 8\sqrt{2}\sqrt{T} \|f_x(x,t)\|_{L_2(D_T)}, \\ A_2(T) = 8 \left(\|p_0(t)\|_{C[0,T]} + \|p_1(t)\|_{C[0,T]} \right). \quad (26)$$

So, we can prove the following theorem.

Theorem. *Let conditions 1-4 be fulfilled and $q \equiv A_2(T) \cdot T \leq 1$ where the numbers $A_2(T)$ is determined by relation (26). Then problem (1)-(3) has a classic solution.*

Proof. Write equation (10) in the form

$$u = \Phi u, \quad (27)$$

where the operator Φ is defined by the right hand side of (10).

Let's consider the operator Φ fixed in the ball $K = K_R \left(\|u\|_{B_{2,T}^{3,3}} \leq R \right)$ where R is such that

$$A_1(T) + A_2(T) \cdot T \leq R,$$

from the space $B_{2,T}^{3,3}$, it is seen from (25) that under the conditions of the theorem for any $u \in K_R$ it holds the inequality

$$\|\Phi u\|_{B_{2,T}^{3,3}} \leq A_1(T) + A_2(T)TR \leq R \tag{28}$$

and for any $u_1, u_2 \in K_R$ we have

$$\|\Phi u_1 - \Phi u_2\|_{B_{2,T}^{3,3}} \leq A_2(T)T \|u_1 - u_2\|_{B_{2,T}^{3,3}} \tag{29}$$

moreover, by conditions of the theorem, $q < 1$.

It follows from inequalities (28) and (29) that under the conditions of the theorem the operator Φ acts in the ball $K = K_R$ and it is contractive. Therefore, in the ball $K = K_R$ the operator Φ has a unique fixed point $\{U\}$ and this point is the solution of (27).

The function $u(x, t)$ as an element of the space $B_{2,T}^{3,3}$ has continuous derivatives $u_x(x, t)$, $u_{xx}(x, t)$, $u_t(x, t)$, $u_{tx}(x, t)$, $u_{ttx}(x, t)$.

Obviously,

$$|u_{ttxx}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} \left(\sum_{k=1}^{\infty} \left(\lambda_k^3 \|u_k''(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}}, \tag{30}$$

$$|u_{ttt}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} \left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k'''(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}}, \tag{31}$$

$$|u_{tttt}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} \left(\sum_{k=1}^{\infty} \left(\lambda_k \|u_k^{(4)}(t)\|_{C[0,T]} \right)^2 \right)^{\frac{1}{2}}. \tag{32}$$

It follows from (31)-(32) that allowing for (22)-(24) the functions $u_{tt}(x, t)$, $u_{ttxx}(x, t)$, $u_{ttt}(x, t)$, $u_{tttt}(x, t)$ are continuous in D_T .

It is easily verified that equation (1) and conditions (2) and (3) are satisfied in an ordinary sense. So $u(x, t)$ is a classic solution of problem (1)-(3). The theorem is proved.

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