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**ON THE BASIS IN THE SPACE  $L_p(0, 1)$ ,  $1 < p < +\infty$   
OF THE SYSTEM OF EIGEN FUNCTIONS OF  
STURM-LIOUVILLE PROBLEM WITH A  
SPECTRAL PARAMETER IN BOUNDARY  
CONDITIONS**

**Abstract**

*We consider the following spectral problem*

$$-y''(x) = \lambda y(x), \quad x \in (0, 1),$$

$$(a_0\lambda + b_0)y(0) = (c_0\lambda + d_0)y'(0),$$

$$(a_1\lambda + b_1)y(1) = (c_1\lambda + d_1)y'(1),$$

*where  $\lambda$  is a spectral parameter,  $a_i, b_i, c_i, d_i, i = \overline{0, 1}$  are real constants, moreover*

$$\sigma_0 = a_0d_0 - b_0c_0 < 0, \quad \sigma_1 = a_1d_1 - b_1c_1 > 0.$$

*Necessary and sufficient basicity conditions in the space  $L_p(0, 1)$ ,  $1 < p < \infty$  of the system of eigen functions of this problem with two removed functions are found.*

Consider the following boundary value problem

$$-y''(x) + q(x)y(x) = \lambda y(x), \quad x \in (0, 1), \quad (1)$$

$$(a_0\lambda + b_0)y(0) = (c_0\lambda + d_0)y'(0), \quad (2)$$

$$(a_1\lambda + b_1)y(1) = (c_1\lambda + d_1)y'(1), \quad (3)$$

where  $\lambda$  is a spectral parameter,  $q(x)$  is a real continuous function on  $[0, 1]$ ,  $a_i, b_i, c_i, d_i, i = \overline{0, 1}$  are real constants, moreover

$$\sigma_0 = a_0d_0 - b_0c_0 < 0, \quad \sigma_1 = a_1d_1 - b_1c_1 > 0. \quad (4)$$

The problem of the form (1)-(3) arises, for example, by separating variables in a dynamic boundary value problem describing small torsional vibrations of a bar with both pulley stiffened ends. A more complete information on the physical sense of the problems of type (1)-(3) may be found in [1] and [2].

In the paper [3] the complete description of general characteristics of arrangement of eigenvalues on a real axis is given, vibrational properties of eigenfunctions are studied, asymptotic formulae for eigenvalues and eigenfunctions of problem (1)-(3) are obtained. The basis properties of eigenfunctions where it is established that the system of eigenfunctions of this problem after removing two arbitrary functions having different parity ordinal number forms a basis in the space  $L_p$ ,  $1 < p < \infty$  is also investigated.

In the paper [4], problem (1)-(3) is reduced to an eigen value problem for a linear operator acting in the Hilbert space  $H = L_2(0, 1) \oplus C^2$ , necessary and sufficient condition of basicity in  $L_p(0, 1)$ ,  $1 < p < \infty$ , of the subsystem of eigen functions of this problem is established. More exactly, it is proved the following theorem.

**Theorem A.** *Let  $r$  and  $l$  be arbitrary fixed entire non-negative numbers. If*

$$\Delta(r, l) = \begin{vmatrix} 1 & 1 \\ c_1 y'_r(1) - a_1 y_r(1) & c_1 y'_l(1) - a_1 y_l(1) \end{vmatrix} \neq 0. \quad (5)$$

then the system of eigen functions  $\{y_k\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) forms a basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ , for  $p = 2$  the Riesz basis, if  $\Delta(r, l) = 0$ , this system is incomplete and not minimal in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ .

The basis properties of the system of root functions in the space  $L_p(0, 1)$ ,  $1 < p < \infty$  of problem (1)-(3) (in special cases) were investigated also in the papers [5] and [6]. In the case  $q \equiv 0$ ,  $b_j = c_j = 0$ ,  $(-1)^{j+1} a_j > 0$ ,  $d_j = 1$ ,  $j = \overline{0, 1}$ , (therewith  $\sigma_0 < 0$ ,  $\sigma_1 > 0$ ) in [5] it was proved that if  $a_0 \neq a_1$  then the system of eigenfunctions of problem (1)-(3) with two removed arbitrary functions forms a basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ ; if  $a_0 = -a_1$ , then the system of eigenfunctions with two removed arbitrary functions having different parity numbers forms a basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ , if  $a_0 = -a_1$ , then the system of eigenfunction with two arbitrary removed functions having the same order parity, is neither complete nor minimal in  $L_p(0, 1)$ ,  $1 < p < \infty$ . In the case  $q \equiv 0$ ,  $b_j = c_j = 0$ ,  $a_j < 0$ ,  $d_j = 1$ ,  $j = \overline{0, 1}$  (therewith  $\sigma_0 < 0$ ,  $\sigma_1 < 0$ ) necessary and sufficient basicity condition in  $L_p(0, 1)$ ,  $1 < p < \infty$ , of the system of the root functions of problem (1)-(3) with two removed root functions is found in [6].

The present paper is devoted to the investigation of basis properties of the subsystem of eigenfunctions of problem (1)-(3) for  $q \equiv 0$ .

Note that the solution of equation (1) satisfying the initial conditions

$$y(0, \lambda) = c_0 \lambda + d_0, \quad y'(0, \lambda) = a_0 \lambda + b_0 \quad (6)$$

is of the form

$$y(x, \lambda) = (c_0 \lambda + d_0) \cos \sqrt{\lambda} x + (a_0 \lambda + b_0) \frac{\sin \sqrt{\lambda} x}{\sqrt{\lambda}}. \quad (7)$$

Taking into account boundary condition (3), we get

$$\begin{aligned} & \cot \sqrt{\lambda} \{(a_0 \lambda + b_0)(c_1 \lambda + d_1) - (a_1 \lambda + b_1)(c_0 \lambda + d_0)\} = \\ & = \frac{1}{\sqrt{\lambda}} (a_0 \lambda + b_0)(a_1 \lambda + b_1) + (c_0 \lambda + d_0)(c_1 \lambda + d_1) \sqrt{\lambda}. \end{aligned}$$

Thus, the eigenvalues  $\lambda_0 < \lambda_1 < \dots < \lambda_k < \dots$  of problem (1)-(3) are the roots of the equation

$$\cot \sqrt{\lambda} = \frac{(a_0 \lambda + b_0)(a_1 \lambda + b_1) + (c_0 \lambda + d_0)(c_1 \lambda + d_1) \lambda}{\{(a_0 \lambda + b_0)(c_1 \lambda + d_1) - (a_1 \lambda + b_1)(c_0 \lambda + d_0)\} \sqrt{\lambda}} \quad (8)$$

and by (7) has only the eigenfunctions

$$y_k(x) = (c_0\lambda + d_0) \cos \sqrt{\lambda_k}x + (a_0\lambda + b_0) \frac{\sin \sqrt{\lambda_k}x}{\sqrt{\lambda_k}}, \quad k = 0, 1, \dots$$

Then we have

$$\begin{aligned} y_k(1) &= (c_0\lambda_k + d_0) \cos \sqrt{\lambda_k} + (a_0\lambda_k + b_0) \frac{\sin \sqrt{\lambda_k}}{\sqrt{\lambda_k}} = \\ &= \sin \sqrt{\lambda_k} \left\{ (c_0\lambda_k + d_0) \cot \sqrt{\lambda_k} + \frac{1}{\sqrt{\lambda_k}} (a_0\lambda_k + b_0) \right\} = \\ &= \sin \sqrt{\lambda_k} \left\{ \frac{(c_0\lambda + d_0) \{ (a_0\lambda_k + b_0) (a_1\lambda_k + b_1) + (c_0\lambda_k + d_0) (c_1\lambda_k + d_1) \lambda_k \}}{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \} \sqrt{\lambda}} + \right. \\ &\quad \left. + \frac{(a_0\lambda_k + b_0) \{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda_k + d_0) \}}{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \} \sqrt{\lambda}} \right\} = \\ &= \sin \sqrt{\lambda_k} \frac{(c_1\lambda_k + d_1) \{ (a_0\lambda_k + b_0)^2 + (c_0\lambda_k + d_0)^2 \lambda_k \}}{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \} \sqrt{\lambda_k}}. \end{aligned} \quad (9)$$

Note that

$$\sin \sqrt{\lambda_k} = (-1)^k \frac{1}{(1 + \cot^2 \sqrt{\lambda_k})^{1/2}}.$$

We have

$$\begin{aligned} 1 + \cot^2 \sqrt{\lambda_k} &= 1 + \frac{\{ (a_0\lambda_k + b_0) (a_1\lambda_k + b_1) + (c_0\lambda_k + d_0) (c_1\lambda_k + d_1) \lambda_k \}^2}{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \}^2 \lambda_k} = \\ &= \{ \{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda_k + d_0) \}^2 \lambda_k + \\ &\quad + \{ (a_0\lambda_k + b_0) (a_1\lambda_k + b_1) + (c_0\lambda_k + d_0) (c_1\lambda_k + d_1) \lambda_k \}^2 \} \times \\ &\quad \times \{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda_k + d_0) \}^{-2} \lambda_k^{-1} = \\ &= \left\{ (a_0\lambda_k + b_0)^2 (c_1\lambda_k + d_1)^2 + (a_1\lambda_k + b_1)^2 (c_0\lambda_k + d_0)^2 \right\} \lambda_k + \\ &\quad + \left\{ (a_0\lambda_k + b_0)^2 (a_1\lambda_k + b_1)^2 + (c_0\lambda_k + d_0)^2 (a_1\lambda_k + b_1)^2 \lambda_k \right\} \times \\ &\quad \times \{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda_k + d_0) \}^{-2} \lambda_k^{-1} = \\ &= \frac{\left\{ (a_0\lambda_k + b_0)^2 (c_0\lambda_k + d_0)^2 \lambda_k (a_1\lambda_k + b_1)^2 + (c_1\lambda_k + d_1)^2 \lambda_k \right\}}{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \}^2 \lambda_k}. \end{aligned}$$

Taking into account the last two equalities, from (9) we find

$$\begin{aligned} y_k(1) &= (-1)^k \times \\ &\times \frac{\{ (a_0\lambda_k + b_0) (c_1\lambda_k + d_1) - (a_1\lambda_k + b_1) (c_0\lambda + d_0) \} \sqrt{\lambda}}{\left\{ (a_0\lambda_k + b_0)^2 + (c_0\lambda_k + d_0)^2 \lambda_k \right\}^{1/2} \left\{ (a_1\lambda_k + b_1)^2 + (c_1\lambda_k + d_1)^2 \lambda_k \right\}^{1/2}} \times \end{aligned}$$

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$$\begin{aligned} & \times \frac{(c_1 \lambda_k + d_1) \left\{ (a_0 \lambda_k + b_0)^2 + (c_0 \lambda_k + d_0)^2 \lambda_k \right\}}{\left\{ (a_0 \lambda_k + b_0) (c_1 \lambda_k + d_1) - (a_1 \lambda_k + b_1) (c_0 \lambda_k + d_0) \right\} \sqrt{\lambda_k}} = \\ & = (-1)^k (c_1 \lambda_k + d_1) \left( \frac{(a_0 \lambda_k + b_0)^2 + (c_0 \lambda_k + d_0)^2 \lambda_k}{(a_1 \lambda_k + b_1)^2 + (c_1 \lambda_k + d_1)^2 \lambda_k} \right)^{1/2}. \end{aligned}$$

Thus, we have

$$y_k(1) = (-1)^k (c_1 \lambda_k + d_1) \left( \frac{(a_0 \lambda_k + b_0)^2 + (c_0 \lambda_k + d_0)^2 \lambda_k}{(a_1 \lambda_k + b_1)^2 + (c_1 \lambda_k + d_1)^2 \lambda_k} \right)^{1/2}. \quad (10)$$

Let the following relation be fulfilled

$$a_1 = a_0, \quad b_1 = b_0, \quad c_1 = -c_0, \quad d_1 = -d_0. \quad (11)$$

Then from (10) we get

$$y_k(1) = (-1)^k (c_1 \lambda_k + d_1). \quad (12)$$

If  $(c_j \lambda_r + d_j) (c_j \lambda_l + d_j) \neq 0$ ,  $j = 0, 1$  then by (14) from [4] (see also [3]), equalities (6) and (12) we have

$$\begin{aligned} \Delta(r, l) = \Delta_1(r, l) &= \begin{vmatrix} 1 & 1 \\ c_1 y'_r(1) - a_1 y_r(1) & c_1 y'_l(1) - a_1 y_l(1) \end{vmatrix} = \\ &= \begin{vmatrix} 1 & 1 \\ \frac{b_1 c_1 - a_1 d_1}{c_1 \lambda_r + d_1} y_r(1) & \frac{b_1 c_1 - a_1 d_1}{c_1 \lambda_l + d_1} y_l(1) \end{vmatrix} = \\ &= \begin{vmatrix} 1 & 1 \\ \frac{y_r(1)}{c_1 \lambda_r + d_1} & \frac{y_l(1)}{c_1 \lambda_l + d_1} \end{vmatrix} = \begin{vmatrix} 1 & 1 \\ (-1)^r & (-1)^l \end{vmatrix}. \end{aligned} \quad (13)$$

If  $c_0 \neq 0$ ,  $\lambda_l = -d_0/c_0$ , then  $\lambda_l = -d_1/c_1$ ,  $c_j \lambda_r + d_j \neq 0$ ,  $j = 0, 1$ . Consequently, by (12) we have

$$\Delta(r, l) = \Delta_5(r, l) = \begin{vmatrix} 1 & 1 \\ \frac{y_r(1)}{c_1 \lambda_r + d_1} & -\frac{c_1 y'_l(1)}{\sigma_1} \end{vmatrix} = \begin{vmatrix} 1 & 1 \\ (-1)^r & -\frac{c_1 y'_l(1)}{\sigma_1} \end{vmatrix}. \quad (14)$$

From formula (7) we get

$$\begin{aligned} y_l(1) = y(1, \lambda_l) &= (c_0 \lambda_l + d_0) \cos \sqrt{\lambda_l} + (a_0 \lambda_l + b_0) \frac{\sin \sqrt{\lambda_l}}{\sqrt{\lambda_l}} = \\ &= (a_0 \lambda_l + b_0) \frac{\sin \sqrt{\lambda_l}}{\sqrt{\lambda_l}} = 0. \end{aligned}$$

Taking into account this equality, from formula (7) we find

$$\begin{aligned} y'_l(1) = y'(1, \lambda_l) &= -\sqrt{\lambda_l} (c_0 \lambda_l + d_0) \sin \sqrt{\lambda_l} + (a_0 \lambda_l + b_0) \cos \sqrt{\lambda_l} = \\ &= \left( a_0 \left( -\frac{d_0}{c_0} \right) + b_0 \right) \cos \sqrt{\lambda_l} = -\frac{a_0 d_0 - b_0 c_0}{c_0} \cos \sqrt{\lambda_l} = \end{aligned}$$

$$= -\frac{a_1 d_1 - b_1 c_1}{c_1} \cos \sqrt{\lambda_l} = -\frac{\sigma_1}{c_1} \sqrt{1 - \sin^2 \sqrt{\lambda_l}} = -\frac{(-1)^k \sigma_1}{c_1}.$$

Taking into account this relation in (14), we get

$$\Delta_5(r, l) = \begin{vmatrix} 1 & 1 \\ (-1)^r & (-1)^l \end{vmatrix}. \tag{15}$$

From (13) and (15) it follows that if  $r$  and  $l$  are entire non-negative numbers having the same parity, then  $\Delta_{r,l} = 0$  and by theorem A, the system of eigenfunctions  $\{y_k\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < +\infty$ .

Thus, we proved

**Theorem 1.** *Let  $r$  and  $l$  be entire non-negative numbers having the same parity, and condition (11) be fulfilled. Then the system of eigenfunctions  $\{y_k\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < +\infty$ .*

From this theorem it is seen that the condition of theorem 4 from [3] about that the numbers  $r$  and  $l$  have different parities is essential.

Let  $r$  and  $l$  be entire non-negative numbers having the same parities. Then it holds the equality

$$\begin{aligned} \Delta(r, l) &= \Delta_1(r, l) = \begin{vmatrix} 1 & 1 \\ \frac{y_r(1)}{c_1 \lambda_r + d_1} & \frac{y_l(1)}{c_1 \lambda_l + d_1} \end{vmatrix} = \\ &= (-1)^r \begin{vmatrix} 1 & 1 \\ \left( \frac{(a_0 \lambda_l + b_0)^2 + (c_0 \lambda_l + d_0)^2 \lambda_l}{(a_1 \lambda_l + b_1)^2 + (c_1 \lambda_l + d_1)^2 \lambda_l} \right)^{1/2} & \left( \frac{(a_0 \lambda_r + b_0)^2 + (c_0 \lambda_r + d_0)^2 \lambda_r}{(a_1 \lambda_r + b_1)^2 + (c_1 \lambda_r + d_1)^2 \lambda_r} \right)^{1/2} \end{vmatrix}. \end{aligned} \tag{16}$$

Now, let's consider the function

$$F(\lambda) = \left( \frac{(a_0 \lambda + b_0)^2 + (c_0 \lambda + d_0)^2 \lambda}{(a_1 \lambda + b_1)^2 + (c_1 \lambda + d_1)^2 \lambda} \right)^{1/2}. \tag{17}$$

Let  $c_0^2 + c_1^2 > 0$ . We rewrite the function  $F(\lambda)$  in the form

$$\begin{aligned} F(\lambda) &= \left( \frac{(a_0 \lambda + b_0)^2 + (c_0 \lambda + d_0)^2 \lambda}{(a_1 \lambda + b_1)^2 + (c_1 \lambda + d_1)^2 \lambda} \right)^{1/2} = \\ &= \left( \frac{c_0^2 \lambda^3 + (a_0^2 + 2c_0 d_0) \lambda^2 + (d_0^2 + 2a_0 b_0) \lambda + b_0^2}{c_1^2 \lambda^3 + (a_1^2 + 2c_1 d_1) \lambda^2 + (d_1^2 + 2a_1 b_1) \lambda + b_1^2} \right)^{1/2}. \end{aligned}$$

Hence we have

$$\begin{aligned} F'(\lambda) &= (1/2) (F(\lambda))^{-1/2} \times \\ &\times \{c_1^2 \lambda^3 + (a_1^2 + 2c_1 d_1) \lambda^2 + (d_1^2 + 2a_1 b_1) \lambda + b_1^2\}^{-2} \times \\ &\times \{(3c_0^2 \lambda^2 + 2(a_0^2 + 2c_0 d_0) \lambda + (d_0^2 + 2a_0 b_0)) \times \end{aligned}$$

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$$\begin{aligned}
& \times (c_1^2 \lambda^3 + (a_1^2 + 2c_1 d_1) \lambda^2 + (d_1^2 + 2a_1 b_1) \lambda + b_1^2) - \\
& \times (c_0^2 \lambda^3 + (a_0^2 + 2c_0 d_0) \lambda^2 + (d_0^2 + 2a_0 b_0) \lambda + b_0^2) \times \\
& \times (3c_1^2 \lambda^2 + 2(a_1^2 + 2c_1 d_1) \lambda + (d_1^2 + 2a_1 b_1)) \} = \\
& = (1/2) (F(\lambda))^{-\frac{1}{2}} \{c_1^2 \lambda^3 + (a_1^2 + 2c_1 d_1) \lambda^2 + (d_1^2 + 2a_1 b_1) \lambda + b_1^2\}^{-2} \times \\
& \times \{c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0)\} \lambda^4 + 2\{c_0^2 (d_1^2 + 2a_1 b_1) - \\
& - c_1^2 (d_0^2 + 2a_0 b_0)\} \lambda^3 + \{3(c_0^2 b_1^2 - c_1^2 b_0^2) + (a_0^2 + 2c_0 d_0) (d_1^2 + 2a_1 b_1) - \\
& - (a_1^2 + 2c_1 d_1) (d_0^2 + 2c_0 d_0) (d_0^2 + 2a_0 b_0)\} \lambda^2 + \\
& + 2\{b_1^2 (a_0^2 + 2c_0 d_0) - b_0^2 (a_1^2 + 2c_1 d_1)\} + \\
& + b_1^2 (d_0^2 + 2a_0 b_0) - b_0^2 (d_1^2 + 2a_1 b_1)\}. \tag{18}
\end{aligned}$$

It follows from formula (18) that either

1)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) \neq 0$ ;

or

2)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,  $c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) \neq 0$ ;

or

3)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,  $c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $c_0^2 b_1^2 - c_1^2 b_0^2 = 0$ , then there exists  $\lambda^* \in R$  such that  $F'(\lambda) \neq 0$  for  $\lambda \geq \lambda^*$ , i.e. the function  $F(\lambda)$  is strongly monotone for  $\lambda \geq \lambda^*$ . Consequently, there exists an entire non-negative number  $k^*$  such that for  $r, l \geq k^*$  we'll have  $\Delta(r, l) = \Delta_1(r, l) \neq 0$  from (16) and (17), i.e. condition (5) is fulfilled. Then on the base of theorem A, the system of eigen functions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$ ,  $r, l \geq k^*$  of problem (1)-(3) for  $q \equiv 0$  forms the Riesz basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ , for  $p = 2$  the Riesz basis.

And if

4)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,  $c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $c_0^2 b_1^2 - c_1^2 b_0^2 = 0$ , then  $F'(\lambda) = 0$  for all  $\lambda \in R$ , i.e.  $F(\lambda) \equiv const$ . Consequently, from (16), (17) we'll have  $\Delta(r, l) = \Delta_1(r, l) = 0$ . Then on the base of theorem A the system of eigenfunctions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ .

Now let  $c_0 = c_1 = 0$ . Then from formula (10) we get

$$y_k(1) = (-1)^k d_1 \left( \frac{(a_0 \lambda + b_0)^2 + d_0^2 \lambda_k}{(a_1 \lambda + b_1)^2 + d_1^2 \lambda_k} \right)^{1/2}. \tag{19}$$

Consider the function

$$F(\lambda) = \left( \frac{(a_0 \lambda + b_0)^2 + d_0^2 \lambda}{(a_1 \lambda + b_1)^2 + d_1^2 \lambda} \right)^{1/2}. \tag{20}$$

We rewrite the function  $F(\lambda)$  in the following form:

$$F(\lambda) = \left( \frac{(a_0 \lambda + b_0)^2 + d_0^2 \lambda}{(a_1 \lambda + b_1)^2 + d_1^2 \lambda} \right)^{1/2} = \left( \frac{a_0^2 \lambda^2 + (2a_0 b_0 + d_0^2) \lambda + b_0^2}{a_1^2 \lambda^2 + (2a_1 b_1 + d_1^2) \lambda + b_1^2} \right)^{1/2}. \tag{21}$$

From (21) we get

$$\begin{aligned}
F'(\lambda) &= \frac{1}{2} (F(\lambda))^{-1/2} \{a_1^2 \lambda + (d_1^2 + 2a_1 b_1) \lambda + b_1^2\}^{-2} \\
&\times \{(2a_0^2 \lambda + (d_0^2 + 2a_0 b_0)) (a_1^2 \lambda^2 + (d_1^2 + 2a_1 b_1) \lambda + b_1^2) - \\
&- (a_0^2 \lambda^2 + (d_0^2 + 2a_0 b_0) \lambda + b_0^2) (2a_1^2 \lambda + (d_1^2 + 2a_1 b_1))\} = \\
&= (1/2) (F(\lambda))^{-1/2} \times \{a_1^2 \lambda + (d_1^2 + 2a_1 b_1) \lambda + b_1^2\}^{-2} \times \\
&\quad \times \{a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0)\} \lambda^2 + \\
&\quad + 2 (a_0^2 b_1^2 - a_1^2 b_0^2) \lambda + \{b_1^2 (d_0^2 + 2a_0 b_0) - b_0^2 (d_1^2 + 2a_1 b_1)\}. \quad (22)
\end{aligned}$$

From (22) it follows that either

1)  $a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0) \neq 0$ ;

or

2)  $a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $a_0^2 b_1^2 - a_1^2 b_0^2 \neq 0$ , then there exists  $\lambda^{**} \in R$  such that  $F'(\lambda) \neq 0$  for  $\lambda \geq \lambda^{**}$ , i.e. the function  $F(\lambda)$  is a strictly monotone for  $\lambda \geq \lambda^{**}$ . Consequently, there exists an entire non-negative number  $k^{**}$  such that for  $r, l \geq k^{**}$  we'll have  $\Delta(r, l) = \Delta_1(r, l) \neq 0$ , from (16), (17). Then on the base of theorem A the system of eigen functions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$ ,  $r, l \geq k^{**}$  of problem (1)-(3) for  $q \equiv 0$  forms a basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ , for  $p = 2$  the Riesz basis.

And if

3)  $a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $a_0^2 b_1^2 - a_1^2 b_0^2 = 0$ , then  $F'(\lambda) = 0$  for all  $\lambda \in R$  i.e.  $F'(\lambda) = const$ . Consequently, from (16), (17) we'll have  $\Delta(r, l) = \Delta_1(r, l) = 0$ . Then again on the base of theorem A, the system of eigenfunctions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ .

So, we proved the following

**Theorem 2.** *Let  $r$  and  $l$  be entire non-negative numbers of the same parity. Then there exists such an entire non-negative number  $\bar{k}$  ( $\bar{k} = \max\{k^*, k^{**}\}$ ), that for  $r, l \geq \bar{k}$  the system of eigen functions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  in the cases*

a)  $c_0^2 + c_1^2 > 0$  and either

1)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) \neq 0$ , or

2)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,  $c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) \neq 0$

or

3)  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,  $c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $c_0^2 b_1^2 - c_1^2 b_0^2 \neq 0$ ,

b)  $c_0 = c_1 = 0$  and either

1)  $a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0) \neq 0$ ; or

2)  $a_0^2 (d_1^2 + 2a_1 b_1) - a_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $a_0^2 b_1^2 - a_1^2 b_0^2 \neq 0$  forms a basis in the spaces  $L_p(0, 1)$ ,  $1 < p < \infty$ , for  $p = 2$  the Riesz basis. In the cases

c)  $c_0^2 + c_1^2 > 0$  and  $c_0^2 (a_1^2 + 2c_1 d_1) - c_1^2 (a_0^2 + 2c_0 d_0) = 0$ ,

$c_0^2 (d_1^2 + 2a_1 b_1) - c_1^2 (d_0^2 + 2a_0 b_0) = 0$ ,  $c_0^2 b_1^2 - c_1^2 b_0^2 = 0$ ;

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d)  $c_0 = c_1 = 0$  and  $a_0^2(d_1^2 + 2a_1b_1) - a_1^2(d_0^2 + 2a_0b_0) = 0$ ,  $a_0^2b_1^2 - a_1^2b_0^2 = 0$ , the system of eigenfunctions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$  is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < +\infty$ .

Note that in the case  $c_0 = c_1 = 0$  and  $b_0 = b_1 = 0$

$$F(\lambda) = \left( \frac{(a_0^2\lambda + d_0^2)}{(a_1^2\lambda + d_1^2)} \right)^{1/2},$$

$$F'(\lambda) = \frac{1}{2} \left( \frac{(a_1^2\lambda + d_1^2)}{(a_0^2\lambda + d_0^2)} \right)^{1/2} \frac{(a_0^2d_1^2 - a_1^2d_0^2)\lambda^2}{(a_1^2\lambda + d_1^2)^2}.$$

Hence it follows that if  $a_0^2d_1^2 - a_1^2d_0^2 \neq 0$ , then  $F'(\lambda) \neq 0$  for  $\lambda \in R \setminus \{0\}$  and consequently the function  $F(\lambda)$  is strongly monotone in  $R$ . From (16) and (17) we'll have  $\Delta(r, l) = \Delta_5(r, l) \neq 0$ . Then on the base of theorem A, the system of eigenfunctions  $\{y_k(x)\}_{k=0, k \neq r, l}^\infty$  of problem (1)-(3) for  $q \equiv 0$  forms a basis in the space  $L_p(0, 1)$ ,  $1 < p < \infty$ , for  $p = 2$  the Riesz basis. But if  $a_0^2d_1^2 - a_1^2d_0^2 = 0$ , then this system is neither complete nor minimal in the space  $L_p(0, 1)$ ,  $1 < p < +\infty$ . Recall that in the case  $d_0 = -1$ ,  $d_1 = 1$  this result was obtained by N.Yu. Kapustin [5].

## References

- [1]. Tikhonov A.N., Samarsky A.A. *Equations of mathematical physics*. M.: Nauka 1972 (Russian).
- [2]. Vladimirov V.S. *Collection of problems on mathematical physics equations*. // M.Nauka, 1982 (Russian).
- [3]. Kerimov N.B., Poladov R.G. *Basis properties of the system of eigen functions of Sturm-Liouville problem with a spectral parameter in boundary conditions*. Dokl. RAN, 2012, vol. 442, No 1, pp. 583-586 (Russian).
- [4]. Aliyev Z.S., Poladov R.G. *Basis properties of eigen functions of Sturm-Liouville problem with a spectral parameter in boundary conditions* // Vestnik Bakinskogo Universiteta, ser. fiz. mat Nauka 2012, No 1, pp. 55-61 (Russian).
- [5]. Kapustin N. Yu. *On spectral problem from mathematical model of the process of vibrations of end pulley bar* // Diff. uravn, 2005, vol. 41, No 10, pp. 1413-1315 (Russian).
- [6]. Aliyev Z.S. *Basis properties of the root functions of a spectral problem with a spectral parameter in boundary conditions* // Dokl. RAN, 2010, vol. 443, No 5, pp. 583-586 (Russian).

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